# Alkali Metal Cation $\pi$ -Interactions in Metalated and Nonmetalated Acetylenes: $\pi$ -Bonded Lithiums in the X-ray Crystal Structures of [Li-C=C-SiMe<sub>2</sub>-C<sub>6</sub>H<sub>4</sub>-OMe]<sub>6</sub> and [Li-O-CMe<sub>2</sub>-C=C-H]<sub>6</sub> and Computational Studies

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**Abstract:** The X-ray crystal structure of  $[\text{Li}-\text{C}=\text{C}-\text{SiMe}_2-\text{C}_6\text{H}_4-\text{OMe}]_6$  (14)<sub>6</sub> features nearly symmetric  $\pi$ -interactions between the lithium ions and the acetylide anions  $(\text{Li}_1-\text{C}_\beta = 2.353(9) \text{ Å}, \text{Li}_1-\text{C}_\alpha = 2.292(9) \text{ Å})$ . These  $\pi$ -contacts are facilitated by the chelating *o*-anisyl methoxy groups  $(\text{Li}_1-\text{C}_\alpha-\text{C}_\beta = 77.6(4)^\circ, \text{Li}_1-\text{O}_1 = 2.169(9)$  Å). The Li-C<sub> $\alpha$ </sub> distances in the  $(\text{LiC}_{\alpha})_6$  core of (14)<sub>6</sub> differ significantly  $(\text{Li}_{1\alpha}-\text{C}_\alpha = 2.132(9) \text{ Å}, \text{Li}_{1b}-\text{C}_\alpha = 2.205(11) \text{ Å})$ . This Li-C<sub> $\alpha$ </sub> distance differentiation is unique in organolithium hexamers, and is due to Li(C=C-R) "side-on- $\pi$ " and "end-on- $\sigma$ " contacts, as is shown computationally in H-C=C-Li(LiH)<sub>2</sub> (20). A second X-ray crystal structure,  $[\text{Li}-\text{O}-\text{CMe}_2-\text{C}=\text{C}-\text{H}]_6$  (22)<sub>6</sub>, reveals electrostatic  $\pi$ -interactions between the lithiums in the (LiO)<sub>6</sub> core and the nonmetalated acetylene groups (Li<sub>1</sub>-C<sub>2</sub> = 2.443(5) Å, Li<sub>1</sub>-C<sub>3</sub> = 2.749(6) Å). These Li-C  $\pi$ -contacts shorten upon acetylene lithiation, as is shown computationally in Li-O-CH<sub>2</sub>-C=C-(H/Li) (24-H/Li). Additional computations reveal that the  $\pi$ -interactions in (HC=C)M<sub>2</sub>H (26-Li-Cs) complexes (modelling oligo- and polymeric M-C=C-R) are weak (only 0.7 kcal/mol for Li), but substantial in M<sup>+</sup>(H-C=C-H) (27-Li-Cs) species (20.2 kcal/mol for Li<sup>+</sup>). In 26-Li-Cs, the  $\pi$ -contacts increase the C=C bond lengths slightly (0.005 Å for Li) and lower the C=C stretching frequencies (33 cm<sup>-1</sup> for Li), they polarize charge density from C<sub> $\alpha$ </sub> toward C<sub> $\beta$ </sub> and hence result in counterion-induced charge delocalizations. The degrees of  $\pi$ -interactions both in (26-Li-Cs) and in (27-Li-Cs) decrease with increasing size of the alkali cations.

#### Introduction

In 1976, Apeloig, Schleyer, Binkley, Pople, and Jorgensen discovered computationally that the dilithioacetylene monomer  $(Li_2C_2)$  prefers a double  $\pi$ -bridged over a linear structure:<sup>1a</sup>

Electrostatic interactions are mainly responsible.<sup>1,2a</sup> This paper is concerned with similar  $\pi$ -interactions both in polar alkali metal acetylides ( $\pi$ -M [M–C=C–R]) and in nonmetalated acetylenes ( $\pi$ -M [H–C=C–R]). Such  $\pi$ -interactions are clearly evident in the structures of alkali metal<sup>2</sup> acetylides when the

**Scheme 1.** Polymer Sheet Structure of the Alkali Metal Acetylides **1-Na-Rb** and **2-Na-Cs**<sup>*a*</sup>



<sup>*a*</sup> The increasing penetration of M-C=C-R layers with increasing size of M is shown.

M−C<sub>α</sub> and M−C<sub>β</sub> distances are similar,<sup>3</sup> e.g., in the polymer sheet arrangements of M−C≡C−H (**1-Na-Rb**)<sup>4</sup> and of M−C≡C−Me (**2-Na-Cs**)<sup>4-6</sup> (Scheme 1, Table 1).

However, such metal  $\pi$ -contacts are not always significant; e.g., note the large  $M-C_{\beta}$  separations (>3 Å) in the oligometric lithium acetylides: [(t-Bu-C=C-Li)<sub>4</sub>(THF)<sub>4</sub>] (3),<sup>7</sup> [(t-Bu-

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C=C-Li)<sub>12</sub>(THF)<sub>4</sub>],<sup>7</sup> [(Ph-C=C-Li)tmpda]<sub>2</sub> (**4**),<sup>8</sup> and [(Ph-C=C-Li)<sub>4</sub>(tmhda)<sub>4/2</sub>] (**5**)<sup>9</sup> (Table 1). In contrast, the short Be- $C_{\beta}$  distances indicate  $\pi$ -interactions in [(Me-C=C)<sub>2</sub>BeNMe<sub>3</sub>]<sub>2</sub> (**6**) (Table 1).<sup>10</sup>



Moreover, the alkali cations in the heterometallic magnesiates  $Li_2[(Ph-C\equiv C)_3Mg(tmeda)]_2$  (7),<sup>11</sup> Na<sub>2</sub>[(t-Bu-C\equiv C)<sub>3</sub>Mg(tmeda)]<sub>2</sub> (8),<sup>12</sup> and Na<sub>2</sub>[(t-Bu-C\equiv C)<sub>3</sub>Mg(pmdta)]<sub>2</sub> (9)<sup>12</sup> connect the acetylene moieties of the [Mg(C≡C-R)<sub>3</sub>]<sup>-</sup> fragments through  $\pi$ -contacts (Table 1). Analogous structures as in 7 to 9 result from replacement of M<sup>+</sup> by EtMg<sup>+</sup> in (Et)(Ph-C≡C)<sub>3</sub>(Mg)<sub>2</sub> (tmeda)]<sub>2</sub>(C<sub>6</sub>H<sub>6</sub>) (10).<sup>13</sup> Similarly, the lithiums in [Me<sub>3</sub>SiC (C≡C-t-Bu)<sub>2</sub>Li]<sub>2</sub> (11)<sup>2g</sup> as well as the magnesium ion in [(C<sub>3</sub>HMe<sub>4</sub>)<sub>2</sub>Ti(C≡C-SiMe<sub>3</sub>)<sub>2</sub>][Mg(THF)CI] (12)<sup>14</sup> are located between the arms of "tweezers" formed by the acetylene groups (Table 1).



How do electrostatic  $\pi$ -interactions affect the electronic structures in metal acetylides? The penetration of the alkali cations into the acetylide layers in **1-Na-Rb** and **2-Na-Cs** increases as the counterions become larger (Scheme 1).<sup>4</sup> The IR  $\nu$ -C=C stretching frequencies of **1-Na-Cs**, **2-Li-Cs**, and **13-**

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**Li-Cs** decrease with increasing size of the alkali metal cations (Table 2).<sup>15</sup> Negative charge delocalization from  $C_{\alpha}$  to  $C_{\beta}$  is indicated both by the increased metal cation/acetylide interactions upon increased M–C=C–R layer penetration (see Scheme 1)<sup>4</sup> and by the lower  $\nu$ -C=C frequencies.<sup>15</sup> However, an increase in ion size also gives rise to lower  $\nu$ -C=C frequencies (see below).<sup>15</sup> How can the effect of  $\pi$ -coordination be differentiated?

Alkali metal  $\pi$ -bonding to benzene ligands has been investigated extensively owing to its important role in biological ion channels.<sup>16</sup> As "lithium-bonded"<sup>17</sup> cyclopropanes emphasize the analogy to hydrogen-bonded cyclopropanes,<sup>17d,18</sup>  $\pi$ -"lithium-bonded" acetylenes stress the analogy to  $\pi$ -hydrogen-bonded acetylenes.<sup>18a,19</sup> The Li<sup>+20</sup> and LiH<sup>21</sup>  $\pi$ -association energies of acetylene are appreciable and are even larger when the acetylenes are metalated.<sup>22</sup> Notwithstanding, Li<sup>+</sup> $\pi$ -bonding has not been observed experimentally in X-ray crystal structures of homogeneous lithium acetylenes or in compounds with nonmetalated acetylene groups.<sup>2</sup>

For an assessment of electrostatic metal acetylene  $\pi$ -interactions,<sup>23</sup> we now report the X-ray crystal structures of lithiated (Li-C=C-SiMe<sub>2</sub>-C<sub>6</sub>H<sub>4</sub>-OMe) and of nonlithiated (Li-O-CMe<sub>2</sub>-C=C-H) acetylene moieties. Both exhibit  $\pi$ -bonded Li ions. High level computations reveal the structural, the energetic, and the  $\omega$ -C=C vibrational consequences of alkali cation  $\pi$ -interactions in related metal acetylene models and assess the electronic effects of  $\pi$ -coordination.

#### **Results and Discussion**

Syntheses and X-Ray Crystal Structures of Homo-Lithium Acetylenes Featuring Electrostatic  $\pi$ -Interactions. Why is  $\pi$ -coordination not apparent in the structures of the oligomeric lithium acetylides 3, 4, and 5? This may be due to (a) insufficient energy gain upon  $\pi$ -bridging (see below), (b) the lower tendency of the smaller alkali metals to undergo multihapto coordination,<sup>2b,c,f</sup> and (c) the competition between the substituents on the acetylene moieties and the lithium coordinating solvent (Scheme 2a).

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**Table 1.** M-C<sub> $\alpha$ </sub>, M-C<sub> $\beta$ </sub>, and C<sub> $\alpha$ </sub>-C<sub> $\beta$ </sub> Distances (Å) in Alkali and Alkaline Earth Metal Acetylenes

	$M=C_{\alpha}{}^{a}$	$M-C_{\beta}$	$C_{\alpha}-C_{\beta}$
$M-C=C-H^b$			
M = Na, 1-Na	2.49(5)/2.7	3.0	1.17(6)
M = K, 1-K	2.87/3.0	3.3	1.2
M = Rb, 1-Rb	2.98/3.2	3.4	1.2
$M-C \equiv C-Me^b$			
M = Na, 2-Na	2.37(15)/2.7	2.8	1.09(20)
M = K, 2-K	2.55(5)/3.0	3.1	1.19(6)
$(t-Bu-C=C-Li)_4(THF)_4, 3^c$	2.19	>3.1	1.20
$[(Ph-C=C-Li)tmpda]_2, 4^d$	2.13/2.16	3.08	1.24
$[(Ph-C=C-Li)_4(tmhda)_2], 5^e$	2.20	>3.1	2.20
$[(MeC \equiv C)_2 BeNMe_3]_2, 6^{f}$	1.763/2.042	2.538	
$Li_2[(PhC \equiv C)_3Mg(tmeda)]_2, 7^g$	2.32	2.48	1.22
Na <sub>2</sub> [(t-Bu-C=C) <sub>3</sub> Mg(tmeda)] <sub>2</sub> , $8^{h,i}$	2.571(3)	2.974(3)	1.200(4)
$Na_2[(t-Bu-C=C)_3Mg(pmdta)]_2, 9^{h,i}$	2.590(5)	2.900(5)	118.5(7)
$[(Et)(PhC \equiv C)_3(Mg)_2(tmeda)]_2(C_6H_6), 10^{i,j}$	2.265(4)	2.678(4)	1.220(6)
$[Me_3SiC(C = C - t - Bu)_2Li]_2, 11^k$	2.11	2.34	
$[(C_5HMe_4)_2Ti(C \equiv C - Si]Me_3)_2][Mg(THF)Cl], 12^l$	2.269(8)	2.456(9)	1.22(1)
$[Li-C=C-SiMe_2-C_6H_4OMe]_6, (14)_6^m$	2.132(9)/2.205(11)/ 2.292(9)	2.353(9)	1.217(6)

<sup>*a*</sup> The shorter  $\sigma$ -(M–C<sub> $\alpha$ </sub>) and the longer  $\pi$ -(M–C<sub> $\alpha$ </sub>) distances are given. <sup>*b*</sup> Reference 4. <sup>*c*</sup> Reference 7. <sup>*d*</sup> Reference 8. <sup>*e*</sup> Reference 9. <sup>*f*</sup> Reference 10. <sup>*s*</sup> Reference 11. <sup>*h*</sup> Reference 12. <sup>*i*</sup> Selected bond distances. <sup>*j*</sup> Reference 13. <sup>*k*</sup> Reference 2g. <sup>*l*</sup> Reference 14. <sup>*m*</sup> See Figure 1.

**Table 2.** Experimental  $\nu$ -C=C Frequencies (cm<sup>-1</sup>) of Alkali Metal Acetylides (see Scheme 1 for the Structures)

М	$M-C \equiv C-H^a 1,$ <b>1-Li-Cs</b>	$M-C \equiv C-Me^a 2,$ 2-Li-Cs	M−C≡C−Ph <sup><i>a</i></sup> 13, 13-Li-Cs	14
Н	1974	2124	2111	2020 <sup>b</sup>
Li		$2053^{c}$	$2036^{d}$	1980 <sup>c</sup>
Na	$1867^{e}$	$2032^{c}$	$2018^{c}$	
Κ	$1858^{c}$	$2023^{c}$	$2000^{c}$	
Rb	1851 <sup>c</sup>	$2020^{c}$	1990 <sup>c</sup>	
Cs	1838 <sup>c</sup>	$2012^{c}$	$1990^{d}$	

<sup>a</sup> Reference 15. <sup>b</sup> Neat. <sup>c</sup> Nujol mull. <sup>d</sup> KBr disk. <sup>e</sup> Reference 46.

## Scheme 2



Recent computations show that solvation by H<sub>2</sub>O molecules has sufficient energy to overcome the  $\pi$ -interactions between lithiums and the triple bonds in (**1-Li**)<sub>2</sub>.<sup>24</sup>



The goal of this research, to realize a homogenous lithium acetylide exhibiting  $\text{Li}-(\text{C}_{\alpha} \equiv \text{C}_{\beta}) \pi$ -interactions, as in (1-Li)<sub>2</sub>, without competing external solvent interactions, led us to examine lithium (*o*-anisyl)dimethylsilylacetylide,  $\text{Li}-\text{C} \equiv \text{C}-\text{SiMe}_2-\text{C}_6\text{H}_4-\text{OMe}$  (14). The *o*-anisyl methoxy group chelation in 14 should result in short  $\text{Li}-\text{C}_{\beta}$  contacts (Scheme 2b).

Indeed, the X-ray crystal structure of hexameric **14** (crystallographic  $S_6$  symmetry) shows such nearly symmetric  $\pi$ -interactions; note the short Li–C<sub> $\beta$ </sub> distances in Figure 1 and Table 1.



**Figure 1.** X-ray crystal structure of  $[\text{Li}-\text{C}=\text{C}-\text{SiMe}_2-\text{C}_6\text{H}_4-\text{OMe}]_6$ (**14**)<sub>6</sub>. Hydrogen atoms are omitted. Selected distances (Å) and angles (deg):  $C_{\alpha}-C_{\beta}$ , 1.217(6);  $\text{Li}_1-C_{\alpha}$ , 2.292(9);  $\text{Li}_1-C_{\beta}$ , 2.353(9);  $\text{Li}_1-O_1$ , 2.169(9);  $\text{Li}_{1a}-C_{\alpha}$ , 2.132(9);  $\text{Li}_{1b}-C_{\alpha}$ , 2.205(11);  $\text{Li}_1-C_{\alpha}-C_{\beta}$ , 77.6(4);  $\text{Li}_{1a}-C_{\alpha}-C_{\beta}$ , 152.0(5);  $\text{Li}_{1b}-C_{\alpha}-C_{\beta}$ , 127.8(5).

The lithium ions (Li<sub>1</sub>) in (14)<sub>6</sub> are coordinated 5-fold by three  $C_{\alpha}$  carbon atoms, by the oxygen atoms of the *o*-anisyl methoxy groups (Li<sub>1</sub>-O<sub>1</sub>: 2.169(9) Å), and by the  $C_{\beta}$  atoms of the acetylene moieties (Li<sub>1</sub>- $C_{\beta}$ : 2.353(9) Å, Table 1). The  $C_{\alpha}$ - $C_{\beta}$ -Si<sub>1</sub> arrangements (177.8(5)°) are nearly linear. As the  $C_{\alpha} \equiv C_{\beta}$ -R fragments tilt toward Li<sub>1</sub>, the Li<sub>(1,1a,1b)</sub>- $C_{\alpha}$ - $C_{\beta}$  angles differ strongly; Li<sub>1</sub>- $C_{\alpha}$ - $C_{\beta}$  (77.6(4)°) is much smaller than Li<sub>1a</sub>- $C_{\alpha}$ - $C_{\beta}$  (152.0(5)°) and the Li<sub>1b</sub>- $C_{\alpha}$ - $C_{\beta}$  angle is intermediate (127.8(5)°). The  $C_{\alpha} \equiv C_{\beta}$  distances (1.217(6) Å) in (14)<sub>6</sub> are increased relative to acetylene 1 (exp, 1.20 Å;<sup>25</sup> calc, 1.199 Å; see below). The methyl groups (C<sub>5</sub>) on O<sub>1</sub> bend 15° out of the plane of the aryl rings (torsion angle  $C_5$ - $O_1$ - $C_{66}$ - $C_{62} = -165°$ ).

Comparisons of  $(14)_6$  and the organolithum hexamers [(c- $\{C_6H_{11}\}Li)_6$  (15)<sub>6</sub>·(C<sub>6</sub>H<sub>6</sub>)<sub>2</sub>],<sup>26</sup> [c- $\{(Me_2C)_2CH\}CH_2Li]_6$  (16)<sub>6</sub>,<sup>27</sup> [Me<sub>3</sub>SiCH<sub>2</sub>Li]<sub>6</sub> (17)<sub>6</sub>,<sup>28</sup> (n-BuLi)<sub>6</sub> (18)<sub>6</sub>,<sup>29</sup> and (i-PrLi)<sub>6</sub> (19)<sub>6</sub><sup>30</sup> are instructive. The (LiC<sub> $\alpha$ </sub>)<sub>6</sub> cores in all these hexameric

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**Table 3.** Bond Distances (Å) and Angles (deg) of the  $(\text{LiC}_{\alpha})_6$  Cores in Hexameric Organolithium Compounds (Scheme 3)

	$Li_1 - Li_{1a}$	$Li_{1a}-Li_{1b}$	$Li_1 - Li_{1b}$	$Li_1-Li_{1b}-Li_{1a}$	$\alpha^{g}$	$Li_1 - C_{\alpha}$	$Li_{1a}-C_{\alpha}$	$Li_{1b}C_{\alpha}$
$[(c-\{C_6H_{11}\}Li)_6 (15)_6 \cdot (C_6H_6)_2]^a$	2.968	2.397	2.397	74.4	70.3	$2.184^{h}$	$2.184^{h}$	2.300
$[c-{(Me_2C)_2CH}CH_2Li]_6, (16)_6^b$	2.976	2.462	2.462	76.5	72.2	2.159	2.123	2.297
$[Me_3SiCH_2Li]_6, (17)_6^c$	3.18	2.45	2.45	80.9	79.5	$2.20^{h}$	$2.20^{h}$	2.28
$(n-BuLi)_6, (18)_6^d$	2.939	2.429	2.429	74.5	70.3	$2.159^{h}$	$2.159^{h}$	2.270
$(i-PrLi)_6, (19)_6^e$	2.959	2.395	2.395	76.3	72.5	$2.180^{h}$	$2.180^{h}$	2.308
( <b>14</b> ) <sub>6</sub> <sup>f</sup>	3.700	2.794	2.794	82.9	80.5	2.292	2.132	2.205

<sup>*a*</sup> Reference 26. <sup>*b*</sup> Reference 27. <sup>*c*</sup> Reference 28. <sup>*d*</sup> Reference 29. <sup>*e*</sup> Reference 30. <sup>*f*</sup> See Figure 1. <sup>*g*</sup> Back to seat angle  $\alpha$  of the Li<sub>6</sub> chair, Scheme 3b. <sup>*h*</sup> Average values of similar distances, which do not differ in more than 0.04 Å; see reference.



organolithium compounds are formed by two stacked  $(LiC_{\alpha})_3$ rings (Scheme 3a). Folded  $Li_6$  chairs are apparent in the  $(LiC_{\alpha})_6$ units (Scheme 3b). The  $C_{\alpha}$  atoms cap the  $Li_3$  faces, with one long  $(Li_1-Li_{1a})$  and two short  $(Li_1-Li_{1b})$  distances (Scheme 3c, Table 3).

Relative to  $(15)_6$ ,  $(16)_6$ ,  $(17)_6$ ,  $(18)_6$ , and  $(19)_6$ , the lithium acetylide  $(14)_6$  exhibits a Li<sub>3</sub> triangle with unusual long Li–Li distances and a rather flat Li<sub>6</sub> chair (large "back-to-seat" angle  $\alpha$ , Scheme 3b, Table 3). In  $(15)_6$  to  $(19)_6$ , the distances of the  $C_{\alpha}$  caps to Li<sub>1</sub> and Li<sub>1a</sub> are short (Scheme 3c, Table 3). Longer Li<sub>1b</sub>-C<sub> $\alpha$ </sub> distances connect the two stacked (LiC<sub> $\alpha$ </sub>)<sub>3</sub> rings (Scheme 3a, Table 3). In  $(14)_6$ , however, the Li<sub>1</sub>-C<sub> $\alpha$ </sub> distances are significantly longer than the Li<sub>1b</sub>-C<sub> $\alpha$ </sub> bonds between the (LiC<sub> $\alpha$ </sub>)<sub>3</sub> subunits (Table 3).

Our computational model for  $(14)_6$ , H–C=C–Li(LiH)<sub>2</sub> (20) (Figure 2), reveals that energy gain upon bending of the  $C_{\alpha} \equiv C_{\beta}$ –H fragment is low (1.36 kcal/mol in 20, Figure 2) but that the  $\pi$ -coordination, which results from the tilt of the  $C_{\alpha} \equiv C_{\beta}$ –R units, gives rise to the unusual Li–C<sub> $\alpha$ </sub> distance differentiations in (14)<sub>6</sub>: Shorter  $\sigma$ -Li–C<sub> $\alpha$ </sub> (2.027 Å) and longer  $\pi$ -Li–C<sub> $\alpha$ </sub> (2.083 Å) distances are apparent in 20-C<sub>s</sub> relative to 20-C<sub>2v</sub> Li–C<sub> $\alpha$ </sub> (2.040 Å, Figure 2). This is due to short "end-on- $\sigma$ " and long "side-on- $\pi$ " C<sub> $\alpha$ </sub> contacts of the Li ions, which coordinate to the  $\sigma$ - and the  $\pi$ -regions of the acetylide ions (Scheme 4, Figure 3).<sup>2c,b</sup>

Deprotonation or metalation increases the affinity of acetylene groups toward metal ion  $\pi$ -coordination (Figure 4),<sup>22</sup> but the Li<sup>+</sup> and LiH  $\pi$ -interaction energies of nonmetalated acetylene are relatively large (22.2 kcal/mol for  $\pi$ -Li<sup>+</sup>(H–C=C–H); see below).<sup>22,31</sup> The short contacts between lithiums and the carbon atoms in lithium pinacolone enolate [Li–O–C(t-Bu)=CH<sub>2</sub>]<sub>6</sub>, (**21**)<sub>6</sub> document  $\pi$ -interactions with Li ions in the (LiO)<sub>6</sub> cluster (Scheme 5).<sup>32</sup> Are analogous electrostatic  $\pi$ -interactions<sup>33</sup> with (nonmetalated) acetylene moieties possible? To provide an answer, we synthesized and crystallized the lithium acetylene alkoxide Li–O–CMe<sub>2</sub>C=C–H (**22**). The X-ray crystal analysis reveals a hexameric aggregate [Li–O–CMe<sub>2</sub>–C=C–H]<sub>6</sub> (**22**)<sub>6</sub> with crystallographic S<sub>6</sub> symmetry (Figure 5).

The lithium centers Li<sub>1</sub> in  $(22)_6$  exhibit short contacts to the organic moieties (Li<sub>1</sub>-C<sub>1</sub> = 2.687(5) Å, Li<sub>1</sub>-C<sub>2</sub> = 2.443(5)



**Figure 2.** (a)  $H-C \equiv C-\text{Li}(\text{LiH})_2$  ( $C_{2\nu}$ , **20-C**<sub>2v</sub>): B3LYP/6-311+G\*\* optimized geometry, total energy -100.613 35 au; B3LYP/6-31G\* zeropoint energy 20.13 kcal/mol (NIMAG = 1). (b)  $H-C \equiv C-\text{Li}(\text{LiH})_2$  ( $C_s$ , **20-C**<sub>s</sub>): B3LYP/6-311+G\*\* optimized geometry, total energy -100.61590 au; B3LYP/6-31G\* zero-point energy 20.37 kcal/mol (NIMAG = 0);  $\pi$ -coordination energy relative to **20-C**<sub>2v</sub> = 1.36 kcal/mol.

Scheme 3



Å, and  $\text{Li}_1-\text{C}_3 = 2.749(6)$  Å (Scheme 5, Table 4). This suggests similar electrostatic interactions as in (21)<sub>6</sub> and in [Li– O–C(Me)–(c-CHCH<sub>2</sub>CH<sub>2</sub>)<sub>2</sub>]<sub>6</sub> (23)<sub>6</sub> (Scheme 5).<sup>17d</sup> This interpretation also is supported by the tilt of the O<sub>1</sub>–C<sub>1</sub> moieties toward Li<sub>1</sub> (Li<sub>1</sub>–O<sub>1</sub>–C<sub>1</sub> = 105.3(2)°, Li<sub>1a</sub>–O<sub>1</sub>–C<sub>1</sub> = 130.8(2)°, Li<sub>1a</sub>–O<sub>1</sub>–C<sub>1</sub> = 134.2(2)°; Table 4) and the coplanarity of the O<sub>1</sub>–Li<sub>1</sub> and C<sub>1</sub>–C<sub>2</sub> bonds (Li<sub>1</sub>–O<sub>1</sub>–C<sub>1</sub>–C<sub>2</sub> dihedral angle = 10.4°, Figure 5).<sup>34</sup> The differences in the Li<sub>(1,1a,1b)</sub>-O<sub>1</sub> distances are remarkable: Li<sub>1</sub>–O<sub>1</sub> (1.955(5) Å) are longer than Li<sub>1a</sub>–O<sub>1</sub>

<sup>(29)</sup> Kottke, T.; Stalke, D. Angew. Chem. 1993, 105, 619; Angew. Chem., Int. Ed. Engl. 1993, 32, 580.

<sup>(30)</sup> Siemeling, U.; Redecker, T.; Neumann, B.; Stammler, H.-G. J. Am. Chem. Soc. 1994, 116, 5507.

<sup>(31)</sup> For electrostatic potential computations of  $H-C \equiv C-H$  and  $F-C \equiv C-H$ , see: Clark, D. T.; Adams, D. B. *Tetrahedron* **1973**, *29*, 1887.

<sup>(32) (</sup>a) Williard, P. G.; Carpenter, G. P. J. Am. Chem. Soc. **1985**, 107, 3345. (b) Williard, P. G.; Carpenter, G. P. J. Am. Chem. Soc. **1986**, 108, 462.

<sup>(33)</sup> For a discussion of electrostatic Li–C $\pi$ -interactions in lithium aryls see: Ruhlandt-Senge, K.; Ellison, J. J.; Wehmschulte, R. J.; Pauer, F.; Power, P. P. J. Am. Chem. Soc. **1993**, 115, 11353

<sup>(34)</sup> Probably due to steric hindrance, ideal syn-periplanar arrangements (Li<sub>1</sub>-O<sub>1</sub>-C<sub>1</sub>-C<sub>2</sub> = 0°) are avoided.

**Scheme 4.** Differentiation of  $Li_1-C_{\alpha}$  and  $Li_{1a}-C_{\alpha}$  Bonds in the X-ray Crystal Structure of  $(14)_{6^a}$ 



<sup>*a*</sup> The R-C=C tilt results in "side-on- $\pi$ " (Li<sub>1</sub>) in addition to the "end-to- $\sigma$ " (Li<sub>1a</sub>) coordination. See also Figure 3.



**Figure 3.** (a) MO contour plots (RHF/6-31+G\*) of H–C≡C–Li-(LiH)<sub>2</sub> ( $C_{2\nu}$ , **20-C**<sub>2v</sub>) reflecting  $\sigma$ - ("in-plane" HOMO-1) and  $\pi$ - ("inplane" HOMO) components of the Li–C bonds. (b) MO contour plots (RHF/6-31+G\*) of H–C≡C–Li(LiH)<sub>2</sub> ( $C_s$ , **20-C**<sub>s</sub>) reflecting  $\sigma$ - ("inplane" HOMO-1) and  $\pi$ - ("in-plane" HOMO) components of the Li–C bonds.

(1.877(5) Å) or  $Li_{1b}-O_1$  (1.923(5) Å). Like the long  $Li_1-C_{\alpha}$  bonds in (14)<sub>6</sub>, the long  $Li_1-O_1$  distances in (22)<sub>6</sub> are obviously due to the "side on p"-coordinated  $O_1$  lone pairs.

To assess the energetics of Li  $\pi$ -bonding in the X-ray crystal structure of (22)<sub>6</sub>, monomeric Li $-O-CH_2-C\equiv C-H$  models were computed without (24-H, Figure 6a) and with (24-H-coord, Figure 6b) Li (H $-C\equiv C$ )  $\pi$ -contacts (Scheme 6). Both structures are minima. The Li (H $-C\equiv C$ )  $\pi$ -interaction is rather weak, and 24-H-coord is only 0.64 kcal/mol more stable than 24-H-coord (Figure 6, a and b). As in the X-ray crystal structure (22)<sub>6</sub>, the  $\pi$ -interactions in 24-H-coord are clearly apparent from the short Li(C $\equiv C$ ) distances (2.241 and 2.478 Å, Figure 6b). Lithiation of the acetylene moiety in 24-H-coord (Figure 4)

and results in shorter Li(C=C)  $\pi$ -contacts (2.115 and 2.290 Å, Figure 6c). Moreover, no minimum corresponding to **24-H** (with X = Li) could be optimized at B3LYP/6-31G\*; only **24-Li-coord** resulted.

Structural, Energetic, and Vibrational Effects of  $\pi$ -Interactions in Alkali Metal Acetylides. Similar  $C_{\alpha}$ -M and  $C_{\beta}$ -M distances (see above) and lower C=C stretching frequencies are important indicators of  $\pi$ -bonding in polar metal acetylides. Are energetic, structural, and vibrational  $\pi$ -interaction criteria related with the charge distributions in alkali metal acetylides M-C=C-H (1-Li-Cs, Table 5)? The (HC=C)M<sub>2</sub>H complexes correspond to the X-ray crystal structures of 4 and of 6 without (25-Li-Cs,  $C_{2\nu}$ ) and with (26-Li-Cs,  $C_s$ )  $\pi$ -interactions (Table 6). Alkali metal cation  $\pi$ -coordinations are apparent in the cationic acetylene complexes 27-Li-Cs (Table 7).

All  $C_s$  structures **26-Li-Cs** with  $\pi$ -interactions are slightly more stable (0.73 kcal/mol (Li) to 0.07 kcal/mol (Cs)) than their  $C_{2\nu}$  counterparts without  $\pi$ -contacts **25-Li-Cs** (Table 6). The  $C_s-C_{2\nu}$  energy differences,  $\Delta E$ , and hence the degree of  $\pi$ -interaction, decrease with increasing ion sizes (increasing distances r between M<sup>+</sup> and the C=C centers). This also is apparent with the  $\pi$ -coordination energies  $E_{\text{coord}}$  of **27-Li-Cs** (Table 7); these correlate with  $1/r^3$  (Figure 7).<sup>35</sup> Although **25-**Li and **26-Li** have the largest energy difference, **25-Li** is the only minimum among the  $C_{2\nu}$  species, **25-Li-Cs** (Table 6).<sup>36</sup>



The C=C distances elongate as the metal ions become larger in **25-Li** to **25-Cs**, e.g., from 1.225 Å to 1.234 Å, as well as from 1.224 Å in **1-Li** to 1.231 Å ion **1-Cs**, reflecting the  $\sigma$ -effects of the cations in these  $\pi$ -bonding free structures (Tables 5 and 6, Figure 8). The  $\pi$ -interactions in **27-Li** to **27-Cs** (C=C = 1.205 and 1.201 Å) result only in small C=C lengthenings relative to **1** (1.199 Å) (Table 7, Figure 8). Similarly, the  $\pi$ -interactions in **26-Li** to **26-Cs** (C=C = 1.230 and 1.234 Å) elongate the C=C distances only slightly relative to the corresponding **25-Li-Cs** structures. The C=C lengthenings due to  $\pi$ -interactions (**25-Li-Cs** vs **26-Li-Cs**) are the greatest the smaller the cations (Figure 8). That the C<sub> $\alpha$ </sub>=C<sub> $\beta$ </sub> distance is a poor criterion for  $\pi$ -interactions has been noted.<sup>3,37</sup>

The  $\sigma$ -coordinated cations shift the harmonic  $\omega$ -C=C stretching frequencies to lower values as the metals become larger: from **25-Li** (2007 cm<sup>-1</sup>) to **25-Cs** (1941 cm<sup>-1</sup>) and from **1-Li** (2014 cm<sup>-1</sup>) to **1-Cs** (1955 cm<sup>-1</sup>, Tables 5 and 6, Figure 9). The  $\pi$ -interactions in **27-Li** to **27-Cs** ( $\omega$ -C=C = 2034 and 2053 cm<sup>-1</sup>) lower the  $\omega$ -C=C frequencies relative to **1** ( $\omega$ -C=C =

<sup>(35)</sup> Whereas in alkali metal cation cyclopropane edge complexes  $E_{\text{coord}}$  depends on  $1/r^{2.5}$  (ref 17d), for benzene a  $1/r^n$  ( $n \le 2$ ) correlation was found (ref 16b).

<sup>(36)</sup> An earlier computational study using the CCSD/DZP level without diffuse functions describes (HCC)Li<sub>2</sub>H as a C<sub>2v</sub> transition structure: Bolton, E. E.; Laidig, W. D.; Schleyer, P. v. R.; Schaefer, H. F., III *J. Am. Chem. Soc.* **1994**, *116*, 9602.

<sup>(37)</sup> Tecle, B.; Ilsley, H.; Oliver, J. P. Inorg. Chem. 1981, 20, 2335.



**Figure 4.** (a-c) Electrostatic potential maps (RHF/6-31+G\*) of H–C≡C–H ( $D_{\infty h}$ , 1), H–C≡C<sup>-</sup> ( $C_{\infty v}$ , 1-anion), and Li–C≡C–H ( $C_{\infty v}$ , 1-Li). The negative values (kcal/mol) reflect the cation coordination affinities of the systems.



**Figure 5.** X-ray crystal structure of  $[Li-O-CMe_2-C≡C-H]_6$  (22)<sub>6</sub>. The methyl groups C<sub>5</sub> are disordered. Hydrogen atoms are omitted except the acetylenic hydrogen atom H<sub>3</sub>. Selected distances (Å) and angles (deg): C<sub>1</sub>-C<sub>2</sub>, 1.486(4); C<sub>2</sub>-C<sub>3</sub>, 1.172(5); Li<sub>1</sub>-C<sub>1</sub>, 2.687(5); Li<sub>1</sub>-C<sub>2</sub>, 2.443(5); Li<sub>1</sub>-C<sub>3</sub>, 2.749(6); Li<sub>1</sub>-O<sub>1</sub>, 1.955(5); Li<sub>1</sub><sub>a</sub>-O<sub>1</sub>, 1.878(5); Li<sub>1</sub><sub>b</sub>-O<sub>1</sub>, 1.932(5); Li<sub>1</sub>-O<sub>1</sub>-C<sub>1</sub>, 105.3(2); Li<sub>1</sub><sub>a</sub>-O<sub>1</sub>-C<sub>1</sub>, 130.8(2); Li<sub>1</sub><sub>b</sub>-O<sub>1</sub>-C<sub>1</sub>, 134.2(2).

**Scheme 5.** Electrostatic Interactions in Hexameric Lithium Alkoxide Clusters



2062 cm<sup>-1</sup>; Table 7, Figure 9). Similarly, the  $\pi$ -interactions in **26-Li** (1974 cm<sup>-1</sup>) to **26-Cs** (1937 cm<sup>-1</sup>) result in lower  $\omega$ -C=C stretching frequencies than in the corresponding **25-Li-Cs** structures. The relative reduction of  $\omega$ -C=C stretching frequencies (**25-Li-Cs** vs **26-Li-Cs**) is the stronger the smaller the

**Table 4.** Selected Bond Distances (Å) and Angles (deg) for  $[\text{Li}-\text{O}-\text{C}(\text{t}-\text{Bu})=\text{CH}_2]_6$ , (**21**)<sub>6</sub>,<sup>*a*</sup>  $[\text{Li}-\text{O}-\text{CMe}_2-\text{C}=\text{C}-\text{H})]_6$  (**22**)<sub>6</sub><sup>*b*</sup> and  $[\text{Li}-\text{O}-\text{C}(\text{Me})-(\text{c}-\text{CHCH}_2\text{CH}_2)_2]_6$  (**23**)<sub>6</sub><sup>*c*</sup> (Scheme 5)

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$(21)_6^{a,d}$	$(22)_6^b$	$(23)_6^c$
1.976(9)	1.955(5)	1.937(3)
1.869(9)	1.877(5)	1.881(3)
1.954(9)	1.923(5)	1.926(3)
88.0(9)	105.3(2)	105.6(1)
140.0(4)	130.8(2)	132.6(1)
132.9(4)	132.4(2)	135.1(1)
2.349(9)	2.687(5)	2.680(3)
2.420(8), 2.53 <sup>e</sup>	2.443(5)	2.615(3)
	2.749(5)	2.644(3)
	$\begin{array}{r} (21)_{6}^{a,d} \\ \hline (21)_{6}^{a,d} \\ \hline 1.976(9) \\ 1.869(9) \\ 1.954(9) \\ 88.0(9) \\ 140.0(4) \\ 132.9(4) \\ 2.349(9) \\ 2.420(8), 2.53^{e} \end{array}$	$\begin{array}{c cccc} (21)_6^{a,d} & (22)_6^{b} \\ \hline & (21)_6^{a,d} & (22)_6^{b} \\ \hline & 1.976(9) & 1.955(5) \\ 1.869(9) & 1.877(5) \\ 1.954(9) & 1.923(5) \\ 88.0(9) & 105.3(2) \\ 140.0(4) & 130.8(2) \\ 132.9(4) & 132.4(2) \\ 2.349(9) & 2.687(5) \\ 2.420(8), 2.53^e & 2.443(5) \\ & 2.749(5) \\ \end{array}$

<sup>*a*</sup> Reference 32. <sup>*b*</sup> Figure 5. <sup>*c*</sup> Reference 17d. <sup>*d*</sup> One of two similar asymmetric units in the unit cell with approximate  $S_6$  symmetry. <sup>*e*</sup> Average value of the two asymmetric units.

cations and appears as a useful indicator for  $\pi$ -interactions, especially for the smaller cations (Figure 9).

The  $\sigma$ -effects of the cations give rise to increased ( $C_{\beta} \rightarrow C_{\alpha}$ ) charge polarizations with smaller cation sizes in **1-Li-Cs** and in **25-Li-Cs** relative to **1** (see the  $C_{\alpha}$ ,  $C_{\beta}$  charges in Tables 5 and 6 and in Figure 10). The  $\pi$ -interactions in **26-Li** to **26-Cs** compensate for these counterion induced  $C_{\alpha}$ ,  $C_{\beta}$  charge separations and hence give rise to  $C_{\alpha} \rightarrow C_{\beta}$  charge delocalizations relative to the corresponding **25-Li-Cs** structures. These cation induced charge delocalizations increase with decreasing cation sizes and hence with increasing degrees of the  $\pi$ -interactions (Table 6, Figure 10).

## Conclusions

Intramolecular coordination of the o-anisyl methoxy groups in [Li-C=C-SiMe<sub>2</sub>-C<sub>6</sub>H<sub>4</sub>-OMe]<sub>6</sub> (14)<sub>6</sub> eliminates external solvent effects and facilitates lithium  $\pi$ -interactions with the acetylide moieties (Li<sub>1</sub>-C<sub> $\beta$ </sub> = 2.353(9) Å), as is evident in the X-ray crystal structure analysis of  $(14)_6$ . The strong tilt of the C=C-R fragments in  $(14)_6$  gives rise to clearly different Li- $C_{\alpha}$  bond lengths in the (LiC<sub> $\alpha$ </sub>) core (Li<sub>1</sub>-C<sub> $\alpha$ </sub> = 2.292(9) Å, Li<sub>1a</sub>-C<sub> $\alpha$ </sub> = 2.132(9) Å, Li<sub>1b</sub>-C<sub> $\alpha$ </sub> = 2.205(11) Å), due to "end-on- $\sigma$ " and "side-on- $\pi$ " coordinated acetylide moieties. Such distinct differences in Li– $C_{\alpha}$  bond distances are unprecedented in the X-ray crystal structures of hexameric organolithium compounds. Similar  $\pi$ -interactions are evident in the X-ray crystal structure of [Li-O-CMe<sub>2</sub>-C≡C-H]<sub>6</sub> (22)<sub>6</sub> from short distances between the lithium ions in the (LiO)6 cluster and the nonmetalated acetylene moieties (Li<sub>1</sub>-C<sub>2</sub> = 2.443(5) Å, Li<sub>1</sub>-C<sub>3</sub> = 2.749(6) Å). Although the  $\pi$ -contacts are clearly evident structurally, the computational models for the X-ray crystal structures  $(14)_6$ , as well as  $(22)_6$ , H-C=C-Li(LiH)<sub>2</sub> (20), and Li-O-CH<sub>2</sub>-C=C-H (24-H), point to the weak nature of these  $\pi$ -interactions (1.36 and 0.64 kcal/mol, respectively). Further computations



**Figure 6.** (a) Li–O–CH<sub>2</sub>–C=C–H (*Cs* **24-H**) without Li (C=C)  $\pi$ -contacts: B3LYP/6-311+G\*\* optimized geometry, total energy –198.888 99 au; B3LYP/6-31G\* zero-point energy 31.91 kcal/mol (NIMAG = 0). (b) Li–O–CH<sub>2</sub>–C=C–H (*Cs* **24-H-coord**) with Li (C=C)  $\pi$ -contact; B3LYP/6-311+G\*\* optimized geometry, total energy –198.889 92 au; B3LYP/6-31G\* zero-point energy 31.85 kcal/mol (NIMAG = 0);  $\pi$ -coordination energy relative to **24-H** = 0.64 kcal/mol. (c) Li–O–CH<sub>2</sub>–C=C–Li (*Cs* **24-Li-coord**) with Li (C=C)  $\pi$ -contact: B3LYP/6-311+G\*\* optimized geometry, total energy –205.829 44 au; B3LYP/6-31G\* zero-point energy 26.40 kcal/mol (NIMAG = 0).

**Scheme 6.** The Li (C=C)  $\pi$ -Interaction Model for the X-Ray Crystal Structure (**22**)<sub>6</sub> (Figure 6)



**Table 5.** Bond Distances (Å),<sup>*a*</sup> Harmonic Vibrational Frequencies  $\omega$  (cm<sup>-1</sup>),<sup>*b*</sup> and Natural Charges *q* (aq)<sup>*c*</sup> of Alkali Metal Acetylides M-C=C-H

	$M{-}C_{\alpha}$	$C_{\alpha} \equiv C_{\beta}$	<i>ω</i> -C≡C	$q \mathrm{M}$	$q C_{\alpha}$	$q  \mathrm{C}_{\!eta}$	$q ~\mathrm{H}$
$\overline{1(D_{\infty h})}$	1.063	1.199	2062		-0.224	-0.224	+0.223
1-Li $(C_{\infty v})$	1.919	1.224	2014	+0.937	-0.746	-0.386	+0.195
1-Na $(C_{\infty v})$	2.222	1.225	2002	+0.909	-0.668	-0.436	+0.195
1-K $(C_{\infty v})$	2.666	1.229	1970	+0.950	-0.643	-0.493	+0.185
<b>1-Rb</b> $(C_{\infty v})$	2.848	1.230	1964	+0.948	-0.626	-0.506	+0.184
1-Cs $(C_{\infty v})$	3.057	1.231	1955	+0.961	-0.621	-0.521	+0.182
<b>1-anion</b> $(C_{\infty v})$		1.243	1882		-0.461	-0.694	+0.155

<sup>*a*</sup> B3LYP/6-311+G\*\* (C, H), 6-31G (Li, Na), LanL2DZ, ECP (K, Rb, Cs) optimized structures. <sup>*b*</sup> Unscaled B3LYP frequencies. <sup>*c*</sup> Natural Population Analysis of the B3LYP electron densities, ref 45.

show the  $\pi$ -interactions in alkali metal (HC=C)M<sub>2</sub>H (**26-Li-Cs**) complexes to be weakly stabilizing (0.73 kcal/mol for Li) and to decrease with increasing cation sizes (0.07 kcal/mol for Cs). The  $\pi$ -contacts give rise to slightly increased C=C bond

lengths (up to 0.005 Å for Li), to lowered  $\omega$ -C=C frequencies (up to 33 cm<sup>-1</sup> for Li) and to cation-induced charge delocalization, which increase with decreasing cation size (Cs to Li).

#### **Experimental Section**

The experiments were carried out under an argon atmosphere by using standard Schlenk as well as needle/septum techniques. The solvents were freshly distilled from sodium/benzophenone. Anisole and 2-methyl-3-butyn-1-ol (Aldrich) were distilled prior to use. Sodium acetylide as a toluene/mineral oil suspension and n-BuLi were purchased from Acros. A hexane solution of 6Li-enriched n-Bu6Li was prepared as described by Seebach et al.<sup>38</sup> The NMR spectra were measured on a JEOL GX spectrometer and referenced to TMS or THF: 1H, 400 MHz; 13C, 100.6 MHz; 29Si, 79.4 MHz; 6Li, 58.9 MHz. IR spectra were determined neat or as Nujol mulls between NaCl discs on a Perkin-Elmer 1420 spectrometer. Mass spectral data were obtained on a Varian MAT 311A spectrometer and the elemental analyses (C, H) on Heraeus micro automaton. The X-ray crystal data were collected with an Enraf Nonius CAD4-Mach3 diffractometer using the  $\omega$ -scan method (3.0°  $< 2\Theta < 54.0^{\circ}$ ). The structures were solved by direct methods using SHELXS 86. The parameters were refined with all data by full-matrix least-squares on F<sup>2</sup> using SHELXL93 (G. M. Sheldrick, Göttingen, 1993). All nonhydrogen atoms were refined anisotropically; the hydrogen atoms were fixed in idealized positions using a riding model.  $RI = \sum |F_{o} - F_{c}| / \sum F_{o}$  and  $wR2 = \sum w |(F_{o}2 - F_{c}2)^{2}| / \sum (w(F_{o}2)^{2})^{0.5}$ . Further details are available on request from the Director of the Cambridge Crystallographic Data Center, Lensfield Rd, GB-Cambridge CB2 1 EW, by quoting the journal citation.

 $Li-C \equiv C-SiMe_2-C_6H_4-OMe$  (14). A solution of ca. 6.8 g (0.06 mol) o-lithioanisole in THF, TMEDA, and hexane was prepared from 37.5 mL (0.06 mol) of BuLi (1.6 M) in hexane, 7.0 g (0.06 mol) of TMEDA, 6.5 g (0.06 mol) of anisole and subsequent solvation of the precipitate with ca. 10 mL of THF.39 This solution was added dropwise at 0 °C to 7.7 g (0.06 mol) of dichlorodimethylsilane in 150 mL of diethyl ether. The resulting mixture was stirred at room temperature for 6 h and filtered; the volatile components were removed by distillation. The residue was taken up in 150 mL of diethyl ether and cooled to 0 °C; a suspension of 2.9 g (0.06 mol) of sodium acetylide in 20 mL of diethyl ether was added. The mixture was stirred for 6 h at room temperature. Hydrolysis with H2O/NH4Cl, extraction with diethyl ether, drying over Na<sub>2</sub>SO<sub>4</sub>, and distillation yielded (9.8 g, 52 mmol, 87%) o-anisyldimethylsilylacetylene, HC=C-SiMe<sub>2</sub>-C<sub>6</sub>H<sub>4</sub>-**OMe**: bp 60 °C/1.6 mbar; <sup>1</sup>H NMR (CDCl<sub>3</sub>) δ 7.69 (d, C<sub>6</sub>H<sub>4</sub>), 7.35 (t, C<sub>6</sub>H<sub>4</sub>), 6.97 (t, C<sub>6</sub>H<sub>4</sub>), 6.79 (t, C<sub>6</sub>H<sub>4</sub>), 3.76 (s, OCH<sub>3</sub>), 2.49 (s, CCH), 0.44 (s, Si(CH<sub>3</sub>)<sub>2</sub>); <sup>13</sup>C {<sup>1</sup>H} NMR (CDCl<sub>3</sub>) δ 165.14, 136.88, 132.40, 124.62, 121.53, 110.48 ( $C_6H_4$ ), 95.30 ( $C_\alpha$ ), 89.80 ( $C_\beta$ ), 55.92 (O- $CH_3$ ), 0.84 (Si(CH<sub>3</sub>)<sub>2</sub>); <sup>29</sup>Si {<sup>1</sup>H} NMR (CDCl<sub>3</sub>)  $\delta$  -21.02; IR (neat, cm<sup>-1</sup>) 3280 (v C≡C-H); 3070, 3002 (v C-H arene), 2960, 2900, 2840 (v C-H aliphatic), 2020 ( $\nu$  C=C).

Lithiation of 0.28 g (1.5 mmol) of  $HC \equiv C-SiMe_2-C_6H_4-OMe$ with 0.9 mL (1.5 mmol) of n-BuLi (1.6 M) in THF or hexane solution (-20 °C, then 5 min. RT) afforded Li- $C \equiv C-SiMe_2-C_6H_4-OMe$ (14) (0.27 g, 1.4 mmol, 93% yield in hexane): NMR<sup>40</sup> of <sup>6</sup>Li-14  $\delta$  <sup>1</sup>H (THF- $d_8$ , +25 °C) 8.13 (d, C<sub>6</sub>H<sub>4</sub>), 7.24 (t, C<sub>6</sub>H<sub>4</sub>), 6.88 (t, C<sub>6</sub>H<sub>4</sub>), 6.77 (d, C<sub>6</sub>H<sub>4</sub>), 3.74 (s, OCH<sub>3</sub>), 0.31 (s, Si(CH<sub>3</sub>)<sub>2</sub>);  $\delta$  <sup>13</sup>C{<sup>1</sup>H} (THF- $d_8$ , -90 °C) 171.62 (pentuplet, C<sub>α</sub>), 164.91 (C<sub>6</sub>H<sub>4</sub>-OMe), 138.65 (C<sub>6</sub>H<sub>4</sub>), 131.27 (C<sub>6</sub>H<sub>4</sub>), 127.04 (C<sub>6</sub>H<sub>4</sub>-SiMe<sub>2</sub>), 120.25 (C<sub>6</sub>H<sub>4</sub>), 114.23 (s, C<sub>β</sub>), 109.11 (C<sub>6</sub>H<sub>4</sub>), 55.00 (OCH<sub>3</sub>), 0.97 (Si(CH<sub>3</sub>)<sub>2</sub>);  $\delta$  <sup>6</sup>Li (THF- $d_8$ , -95 °C) -0.26 (s);  $\delta$  <sup>29</sup>Si {<sup>1</sup>H} (THF- $d_8$ , 31 °C) -30.55; IR (Nujol, cm<sup>-1</sup>) 3050 ( $\nu$ C-H arene), 1980 ( $\nu$  C≡C); MS (EI, 70 eV, 120 °C) *m/e* 354 [Me-O-C<sub>6</sub>H<sub>4</sub>-SiMe<sub>2</sub>-C≡C-SiMe<sub>2</sub>-C<sub>6</sub>H<sub>4</sub>-OMe]<sup>+</sup>, 339 [Me-O-C<sub>6</sub>H<sub>4</sub>-SiMe<sub>2</sub> C≡C-SiMe-C<sub>6</sub>H<sub>4</sub>-OMe]<sup>+</sup>, 309 [Me-O-C<sub>6</sub>H<sub>4</sub>-SiMe-C≡C-SiMe-C<sub>6</sub>H<sub>4</sub>-OMe]<sup>+</sup>, 309 [Me-O-C<sub>6</sub>H<sub>4</sub>-SiMe-C=C-SiMe-C<sub>6</sub>H<sub>4</sub>-OMe]<sup>+</sup>, 309 [Me-O-C<sub>6</sub>H<sub>4</sub>-SiMe-C=C-SiMe-C<sub>6</sub>H<sub>4</sub>-OMe]<sup>+</sup>, 309 [Me-O-C<sub>6</sub>H<sub>4</sub>-OMe]<sup>+</sup>, 279

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**Table 6.** Energies,<sup>*a*</sup> Bond Distances (Å),<sup>*a*</sup> Harmonic Vibrational Frequencies  $\omega$  (cm<sup>-1</sup>),<sup>*b*</sup> and Natural Charges q (au)<sup>*c*</sup> of (HC=C)M<sub>2</sub>H Complexes without ( $C_{2\nu}$ ) and with ( $C_s$ )  $\pi$ -Interactions

	E, ZPE (NIMAG) <sup><math>d</math></sup>	$\Delta E^{e}$	$M_1 - C_{\alpha}$	$M_{1a}$ - $C_{\alpha}$	$M-C_{\beta}$	$C_{\alpha} \equiv C_{\beta}$	ω-C≡C	$q \mathbf{M}^{f}$	$q C_{\alpha}$	$q~\mathrm{C}_eta$	$q \mathrm{H}$
25-Li (C <sub>2v</sub> )	-92.445 65, 15.79 (0)		2.091	2.091	3.183	1.225	2007	+0.873	-0.775	-0.337	+0.204
<b>26-Li</b> $(C_s)$	-92.446 81, 15.79 (0)	0.73	2.126	2.051	2.446	1.230	1974	+0.874	-0.652	-0.478	+0.213
<b>25-Na</b> $(C_{2v})$	-401.966 01, 14.03 (1, -79)		2.415	2.415	3.484	1.227	1983	+0.878	-0.703	-0.409	+0.199
<b>26-Na</b> $(C_s)$	-401.967 15, 14.10 (0)	0.65	2.428	2.389	2.929	1.230	1968	+0.879	-0.642	-0.477	+0.202
<b>25-K</b> ( $C_{2v}$ )	-133.658 25, 12.95 (1, -44)		2.834	2.834	3.860	1.231	1956	+0.923	-0.658	-0.487	+0.192
<b>26-K</b> ( $C_s$ )	-133.658 69, 13.02 (0)	0.21	2.838	2.836	3.416	1.233	1950	+0.923	-0.621	-0.523	+0.193
<b>25-Rb</b> $(C_{2v})$	-125.102 71, 12.60 (1, -45)		3.025	3.005	4.035	1.232	1948	+0.924	-0.637	-0.508	+0.190
<b>26-Rb</b> $(C_s)$	-125.103 08, 12.67 (0)	0.16	3.031	3.033	3.595	1.232	1942	+0.923	-0.607	-0.537	+0.191
<b>25-Cs</b> $(C_{2v})$	-117.125 51, 12.35 (1, -41)		3.229	3.229	4.217	1.234	1941	+0.937	-0.626	-0.529	+0.188
<b>26-Cs</b> $(C_s)$	-117.125 74, 12.42 (0)	0.07	3.232	3.243	3.835	1.234	1937	+0.936	-0.605	-0.549	+0.188

<sup>*a*</sup> B3LYP/6-311+G\*\* (C, H), 6-31G (Li, Na), LanL2DZ, ECP (K, Rb, Cs) optimized structures. <sup>*b*</sup> Unscaled B3LYP frequencies. <sup>*c*</sup> Natural Population Analysis of the B3LYP electron densities, ref 45. <sup>*d*</sup> Total energies *E* (au), unscaled zero-point energies ZPE (kcal/mol), numbers and values (cm<sup>-1</sup>) of imaginary frequencies in parentheses. <sup>*e*</sup> Relative  $C_s - C_{2v}$  energies  $\Delta E$  (kcal/mol). <sup>*f*</sup> Average value for M<sub>1</sub> and M<sub>1a</sub>.

**Table 7.** The  $\pi$ - Coordination Energies  $E_{\text{coord}}$  (kcal/mol),<sup>*a*</sup> Bond Distances (Å),<sup>*a*</sup> Harmonic Vibrational Frequencies  $\omega$  (cm<sup>-1</sup>),<sup>*b*</sup> and Natural Charges q (au)<sup>*c*</sup> of the Alkali Cation Acetylene Complexes M<sup>+</sup>(H–C=C–H)

	$E_{ m coord}$	$M^+ \pi - (C \equiv C)^d$	$C_{\alpha} \equiv C_{\beta}$	ω−C≡C	$q \mathrm{M}$	<i>q</i> C	$q~{ m H}$
$1(D_{\infty h})$			1.199	2062		-0.223	+0.223
<b>27-Li</b> $(C_{2v})$	20.22	2.254	1.205	2034	+0.978	-0.261	+0.272
<b>27-Na</b> $(C_{2v})$	13.29	2.633	1.204	2039	+0.986	-0.252	+0.259
<b>27-K</b> $(C_{2\nu})$	7.50	3.152	1.202	2048	+0.997	-0.245	+0.246
<b>27-Rb</b> $(C_{2\nu})$	5.73	3.454	1.201	2051	+0.998	-0.241	+0.242
<b>27-Cs</b> $(C_{2v})$	4.53	3.739	1.201	2053	+0.999	-0.238	+0.239

<sup>*a*</sup> B3LYP/6-311+G\*\* (C, H), 6-31G (Li, Na), LanL2DZ, ECP (K, Rb, Cs) optimized structures. <sup>*b*</sup> Unscaled B3LYP frequencies. <sup>*c*</sup> Natural Population Analysis of the B3LYP electron densities, ref 45. <sup>*d*</sup> Distance between the metal and the center of the C=C bond.



**Figure 7.** The  $1/r^3$  dependence of the coordination energy  $E_{\text{coord}}$  in  $M^+$  (H–C=C–H) complexes (**27-Li-Cs**).



Figure 8. The C=C distances of the alkali metal acetylene compounds.

 $[Me-O-C_6H_4-Si-C=C-Si-C_6H_4-O]^+$ . Anal.  $(C_{11}H_{13}O_1Li_1Si_1)$ . Calcd: C, 67.3; H, 6.6. Found: C, 66.8; H, 6.8. Single crystals of **14** were obtained from cooled hexane solutions.

X-ray crystal data for (14)<sub>6</sub>:  $M_r = 196.24$ ; rhombohedric; space group R-3; a = b = 22.577(3) Å, c = 12.774(2) Å; V = 5638.8(14)Å<sup>3</sup>;  $D_{calc} = 1.040$  Mgm<sup>-3</sup>; Z = 18; F(000) = 1872; Mo K $\alpha$  ( $\lambda =$ 







**Figure 10.** The natural charges on  $C_{\alpha}$  and  $C_{\beta}$  in the alkali metal acetylides.

0.71073 Å); T = 293 (2) K; crystal dimensions:  $0.30 \times 0.20 \times 0.20$  mm; total reflections 2137; unique 1966;  $I > 2\sigma(I)$ , 1009; parameters, 128. Final *R* values: R1 = 0.0917 ( $I > 2\sigma(I)$ ) and wR2 = 0.1759 (all data). GOF = 1.148; largest peak (0.171 e Å<sup>-3</sup>) and hole (-0.150 e Å<sup>-3</sup>).

 $Li-O-CMe_2-C \equiv C-H$  (22). A solution of 0.223 g (2.66 mmol) of 2-methyl-3-butyn-2-ol (H-C  $\equiv C-CMe_2-OH$ ) in 10 mL of hexane

was cooled to 0 °C and 1.66 mL (2.66 mmol) of 1.6 M BuLi in hexane were added. The white suspension (isolated: 0.23 g, 2.56 mmol, 96% yield), stirred at room temperature for 5 min, dissolved on warming. Slowly cooling to room temperature afforded colorless crystals. NMR of **22**:  $\delta$  <sup>1</sup>H (CDCl<sub>3</sub>, +25 °C) 2.35 (s, *H*-C=C), 1.44 (s, *CH*<sub>3</sub>);  $\delta$  <sup>13</sup>C{<sup>1</sup>H} (CDCl<sub>3</sub>, +25 °C) 94.57 (C<sub>1</sub>), 67.67 (C<sub>3</sub>), 64.51 (C<sub>2</sub>), 35.06 (*CH*<sub>3</sub>); IR (Nujol, cm<sup>-1</sup>) 3270 ( $\nu$  C-H); MS (EI, 70 eV, 80 °C) *m/e* 457 [M<sub>5</sub>-Li<sup>+</sup>], 399 [457<sup>+</sup> - COMe<sub>2</sub>], 341 [399<sup>+</sup> - COMe<sub>2</sub>]. Anal. (C<sub>5</sub>H<sub>7</sub>Li<sub>1</sub>O<sub>1</sub>). Calcd: C, 66.7; H, 7.8. Found: C, 66.4; H, 7.9. Single crystals of **22** were obtained from cooled hexane solutions.

X-ray crystal data for (22)<sub>6</sub>:  $M_r = 90.05$ ; rhombohedral, obv.; space group R-3; a = b = 10.767(2) Å, c = 26.933(3) Å, V = 2704.0(9) Å<sup>3</sup>;  $D_{calc} = 0.995$  Mgm<sup>-3</sup>; Z = 18; F(000) = 864; Mo K $\alpha$  ( $\lambda = 0.71073$ Å); T = 293 (2) K; crystal dimensions:  $0.30 \times 0.30 \times 0.20$  mm; total reflections 1438; unique 1234;  $I > 2\sigma(I)$ , 700. H<sub>3</sub> was refined independently and anisotropically; the other hydrogen atoms were fixed in idealized positions using a riding model. Final *R* values: R1 =0.0781 ( $I > 2\sigma(I)$ ) and wR2 = 0.1882 (all data). GOF = 1.185; largest peak (0.192 e Å<sup>-3</sup>) and hole (-0.155 e Å<sup>-3</sup>).

#### **Theoretical Section**

All computed structures were optimized with Becke's threeparameter hybrid functional<sup>41</sup> incorporating the Lee-Yang-Parr correlation functional<sup>42</sup> (Becke3LYP) using the gradient techniques implemented in GAUSSIAN94.<sup>43</sup> The 6-311+G\*\* (C, H, Li) and 6-31G (Li, Na) basis sets were used. For K, Rb, and Cs 9-valence electron effective core potentials and the LanL2DZ basis sets [K, (341/ 311); Rb, (341/321); Cs, (341/321)] were employed.<sup>44</sup> The character of the stationary points, the zero-point energy correction, and the harmonic vibration frequencies were obtained from analytical and, for pseudo-potential computations of the K, Rb, and Cs systems, from numerical frequency calculations. All partial charges are based on Natural Population Analysis (NPA)<sup>45</sup> of the Becke3LYP electron density. The electrostatic potentials were evaluated with RHF/6-31+G\* wave functions on optimized B3LYP geometries.

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**Supporting Information Available:** Tables giving crystal data and structure refinement details, atomic coordinates, bond distances and angles, and thermal parameters for **14** and **22** (12 pages). See any current masthead page for ordering and Internet access instructions.

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